Cohomology of effect algebras

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Alice and Bob select one of two binary measurements. Alice's measurements: a, a' with possible outcomes 0, 1 Bob's measurements: b, b' with possible outcomes 0, 1 Probabilities of joint outcomes:

	(0,0)	(0, 1)	(1, 0)	(1, 1)
(a, b)	1/2	0	0	1/2
(a, b')	3/8	1/8	1/8	3/8
(a', b)	3/8	1/8	1/8	3/8
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- This setup is quantum mechanically realizable.
- It is not classically realizable.

Goal: develop systematic techniques to study realizability.

- Bell-type scenarios can be described using topology of measurement covers. (Abramsky, Brandenburger)
- Using topological cohomology theory, one obtains a criterion for classical realizability. (Abramsky, Mansfield, Soares Barbosa)
- Bell-type scenarios can also be described using effect algebras. (Staton, Uijlen)
- Can we define cohomology of effect algebras?

- Effect algebras
- Ochomology
- O Applications to Bell-type scenarios

Key feature of quantum logic: Partiality

- A = "The particle *P* is at position x_0 ."
- B = "The particle P has momentum p_0 ."

Conjunction $A \wedge B$ is not defined in this case

Effect algebras

Definition

An effect algebra consists of:

- A set A
- A partial binary operation \oplus on A
- Constants $0, 1 \in A$
- An orthocomplement operation $(-)^{\perp}: A \rightarrow A$

such that

- $\bullet\,$ The operation \oplus is commutative and associative and has 0 as neutral element
- For every $a \in A$, a^{\perp} is the unique element for which $a \oplus a^{\perp} = 1$
- $0^{\perp} = 1$
- If $a \oplus 1$ is defined, then a = 0

Examples

- The unit interval [0,1] is an effect algebra, with addition as \oplus and $a^{\perp} = 1 a$.
- Let H be a Hilbert space. Then

$$\mathcal{E}f(H) = \{A : H \to H \mid 0 \le A \le I\}$$

forms an effect algebra with the same operations.

- Any Boolean algebra is an effect algebra. a ⊕ b is defined whenever a ∧ b = 0, and in that case a ⊕ b = a ∨ b.
- Similarly, any orthomodular poset is an effect algebra.

Effects represent measurements on a physical system. To each effect algebra we associate a state space representing the corresponding states.

$$\mathsf{St}(A) = \left\{ \sigma : A \to [0,1] \; \left| \begin{array}{c} \sigma(a \oplus b) = \sigma(a) + \sigma(b) \\ \sigma(1) = 1 \end{array} \right. \right\}$$

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Example

The state space of $\mathcal{P}(n)$ consists of $a_1, \ldots, a_n \in [0, 1]$ such that $a_1 + \cdots + a_n = 1$.



$$\mathsf{St}(A) = \left\{ \sigma : A \to [0,1] \; \left| \begin{array}{c} \sigma(a \oplus b) = \sigma(a) + \sigma(b) \\ \sigma(1) = 1 \end{array} \right\} \right\}$$

Example

The state space of $\mathcal{E}f(H)$ is the set of density matrices on H, i.e. all positive $\rho: H \to H$ for which $tr(\rho) = 1$.

$$\mathsf{St}(A) = \left\{ \sigma : A \to [0,1] \; \left| \begin{array}{c} \sigma(a \oplus b) = \sigma(a) + \sigma(b) \\ \sigma(1) = 1 \end{array} \right\} \right\}$$

The state space always forms a compact convex space: if σ, τ are states and $\lambda \in [0, 1]$, then

$$\lambda \sigma + (1 - \lambda) \tau$$

is again a state.

Bell's experiment using effect algebras





Measurements and probabilities can be modeled by effect algebras.

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Theorem (Staton, Uijlen)

 σ classically realizable $\iff \sigma$ factors through E_{CM}

Functors H^0, H^1, H^2, \ldots : **TopSp** \rightarrow **AbGrp**



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In effect algebra theory:

Functors H^0, H^1, H^2, \ldots : EffAlg \rightarrow AbGrp

 $H^n(A)$ provides information about states and state extensions.



We modify Connes' definition of cyclic cohomology.

1 Tests on an effect algebra *A*:

$$T_n(A) = \{(a_0,\ldots,a_n) \mid a_0 \oplus \cdots \oplus a_n = 1\}$$

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• Tests on an effect algebra A:

$$T_n(A) = \{(a_0,\ldots,a_n) \mid a_0 \oplus \cdots \oplus a_n = 1\}$$

Operations on tests:

$$d_i: \quad T_n(A) \to T_{n-1}(A)$$

$$d_0: (a_0, \dots, a_n) \mapsto (a_0 \oplus a_1, a_2, \dots, a_n)$$

$$d_1: (a_0, \dots, a_n) \mapsto (a_0, a_1 \oplus a_2, a_3, \dots, a_n)$$

$$\vdots$$

$$d_n: (a_0, \dots, a_n) \mapsto (a_n \oplus a_0, a_1, \dots, a_n)$$

$$T_n(A) = \{(a_0,\ldots,a_n) \mid a_0 \oplus \cdots \oplus a_n = 1\}$$

Occycles:

$$\mathcal{C}^{n}(A) = \left\{ \varphi : T_{n}(A) \to \mathbb{R} \mid \begin{array}{c} \varphi(a_{n}, a_{0}, \dots, a_{n-1}) \\ = (-1)^{n} \varphi(a_{0}, \dots, a_{n}) \end{array} \right\}$$

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Example

For n = 1:

$$\varphi(a, b) = -\varphi(b, a)$$

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Example

For n = 1:

$$\varphi(\mathbf{a}, \mathbf{a}^{\perp}) = -\varphi(\mathbf{a}^{\perp}, \mathbf{a})$$

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Operations on cocycles:

$$d^{i}: C^{n-1}(A) \to C^{n}(A)$$
$$d^{i}\varphi = \left(T_{n}(A) \xrightarrow{d_{i}} T_{n-1}(A) \xrightarrow{\varphi} \mathbb{R}\right)$$

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$$\mathcal{C}^n(A) = \left\{ \varphi : T_n(A) \to \mathbb{R} \mid \begin{array}{l} \varphi(a_n, a_0, \dots, a_{n-1}) \\ = (-1)^n \varphi(a_0, \dots, a_n) \end{array} \right\}$$

Ochain complex:

$$\mathcal{C}^{0}(A) \xrightarrow[-d^{1}]{-} d^{0} \xrightarrow[]{\rightarrow} \mathcal{C}^{1}(A) \xrightarrow[-d^{1}]{-} d^{0} \xrightarrow[]{\rightarrow} \mathcal{C}^{2}(A) \xrightarrow[-d^{2}]{-} d^{0} \xrightarrow[]{\rightarrow} \cdots$$

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$$\xrightarrow{-d^{2} \rightarrow} \overset{-d^{2} \rightarrow}{-d^{3} \rightarrow} \cdots$$
$$\delta^{n} = d^{0} - d^{1} + d^{2} - \cdots \pm d^{n}$$

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6 Cochain complex:

$$\mathcal{C}^{0}(A) \xrightarrow{\delta^{1}} \mathcal{C}^{1}(A) \xrightarrow{\delta^{2}} \mathcal{C}^{2}(A) \xrightarrow{\delta^{3}} \cdots$$

$$\delta^n = d^0 - d^1 + d^2 - \dots \pm d^n$$

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•
$$H^n(A) = \ker(\delta^{n+1}) / \operatorname{im}(\delta^n)$$

Examples

• Cohomology of the unit interval [0, 1]:

$$H^0([0,1]) = \mathbb{R}$$

 $H^n([0,1]) = 0$ for $n > 0$

• Cohomology of the Boolean algebra $\mathcal{P}(m)$:

$$H^n(\mathcal{P}(m)) = \mathbb{R}^{\binom{m-1}{n}}$$

The first cohomology group is related to the state space. The state space is always a compact convex space. Every convex space can be embedded in an \mathbb{R} -vector space. In fact, there is a smallest vector space in which it embeds:



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Theorem

Let A be a finite effect algebra that has enough states. Then $H^{1}(A)$ is the smallest vector space in which St(A) can be embedded:



"A has enough states" means: if $\sigma(a) = \sigma(b)$ for all states σ , then a = b. Let $\sigma: A \to [0,1]$ be a state on an effect algebra A, for instance the Bell state. Then:

$$\sigma \text{ classically realizable} \longleftrightarrow \sigma \text{ factors through} \\ a \text{ Boolean algebra } B$$



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$$\begin{array}{ccc} \mathsf{St}(A) & \stackrel{i}{\longrightarrow} & H^1(A) & \stackrel{\partial}{\longrightarrow} & H^2(B,A) \\ \sigma & \longmapsto & \partial(i(\sigma)) \end{array}$$

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How to get rid of false positives?

$$\mathcal{C}^{n}(A) = \left\{ \varphi : T_{n}(A) \to \mathbb{R} \mid \begin{array}{c} \varphi(a_{n}, a_{0}, \dots, a_{n-1}) \\ = (-1)^{n} \varphi(a_{0}, \dots, a_{n}) \end{array} \right\}$$

$$\mathcal{C}^{n}_{\leq}(A) = \left\{ \left(\begin{array}{c} \varphi : T_{n}(A) \to \mathbb{R}, \\ \psi : T_{n-1}(A) \to \mathbb{R} \end{array} \right) \middle| \varphi \geq \delta \psi \right\}$$

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- Effect algebras can be used to model contextuality scenarios.
- Cohomology of effect algebras is relatively easy to compute, and contains information about states and classical realizability.
- Order cohomology provides a criterion for classical realizability without false positives, but is more difficult to compute.