

An operational resource theory of purity

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Introduction

Resource-theoretic approach to the **foundations of thermodynamics**.

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Free operations give a **preorder** on resources: *A more valuable than B* if $A \xrightarrow{\text{free}} B$.

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Free operations of the form (**basic noisy operations**)

$$\mathcal{N}(\rho_A) = \text{tr}_E \left[U_{AE} (\rho_A \otimes \chi_E) U_{AE}^\dagger \right]$$

and its topological closure [Shor, Gour et al.].

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Toy model for thermal interaction!

Is it specific to QM?

Contents

- 1 Operational purity
- 2 Other definitions of purity

Section 1

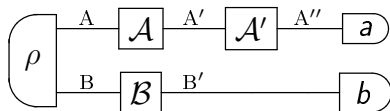
Operational purity

OPTs

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Reversible transformations

A (deterministic) transformation $\mathcal{U} : A \rightarrow B$ is **reversible** if there exists $\mathcal{U}^{-1} : A \rightarrow B$ such that $\mathcal{U}^{-1}\mathcal{U} = \mathcal{I}_A$, and $\mathcal{U}\mathcal{U}^{-1} = \mathcal{I}_B$.

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In general we can't follow the approaches relying on the invariant state.

Purity via RaRe channels

[Chiribella & Scandolo '15b]

Free operations (**RaRe channels**) of the form [Uhlmann]

$$\mathcal{R} = \sum_i p_i \mathcal{U}_i,$$

with \mathcal{U}_i 's *reversible* channels and $\{p_i\}$ *probability distribution*.

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- **No reference to the invariant state!**

Section 2

Other definitions of purity

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Exploring other ways to define purity with the invariant state as free.

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with some thermodynamic flavour!



Purity Preservation

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- Without Purity Preservation, information could be lost simply by composing devices!

Causality

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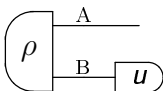
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$$\rho_A := \text{Tr}_B \rho_{AB} =$$


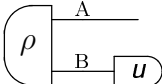
The diagram shows a large rounded rectangle labeled ρ representing a bipartite state. Two horizontal lines extend from the right side of the rectangle, labeled 'A' (top) and 'B' (bottom). The line for system 'B' is connected to a smaller rounded rectangle labeled 'u', representing a deterministic effect.

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Important in thermodynamics: we need to consider subsystems of larger systems!

Purification [Chiribella et al. '10]

- ① Every state ρ_A can be **purified**: there exists a **pure** state Ψ_{AB} such that

$$\rho \text{---} A = \Psi \begin{matrix} \text{---} A \\ \text{---} B \end{matrix} \text{---} U$$

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$$\rho_A \text{---} A = \left(\Psi \begin{array}{l} A \\ B \end{array} \right) \text{---} U$$

- 2 Different purifications of the same state differ by a reversible transformation \mathcal{U} on the **purifying system**:

$$\left(\Psi \begin{array}{l} A \\ B \end{array} \right) \text{---} U = \left(\Psi' \begin{array}{l} A \\ B \end{array} \right) \text{---} U \implies \left(\Psi \begin{array}{l} A \\ B \end{array} \right) = \left(\Psi' \begin{array}{l} A \\ B \end{array} \right) \text{---} \mathcal{U} \text{---} B$$

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It also guarantees the existence of a unique invariant state χ for every system!

Pure Sharpness

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- We can think of a as part of a yes/no test to check an *elementary property* of the system.
- Pure Sharpness guarantees that every system has an elementary property.

Permutability and Strong Symmetry

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Under the previous axioms, Permutability implies Strong Symmetry.

Real, complex, and fermionic QM (+ possible higher-order interference variants) satisfy all these axioms.

Consequences of the axioms

- 1 For every non-trivial system A , an integer $d_A \geq 2$ (**dimension**) gives the size of **all** *maximal* sets of perfectly distinguishable pure states.

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We strongly think that Permutability/Strong Symmetry isn't necessary in fact (work in progress...).

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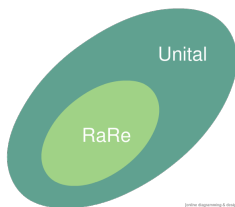
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RaRe channels are **unital**, but **not all** unital channels are RaRe [Landau & Streater]!

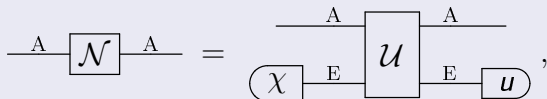


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Noisy operations

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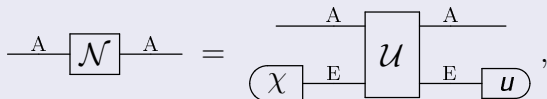


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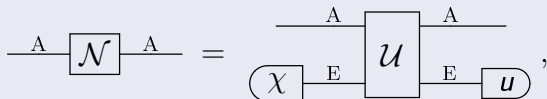


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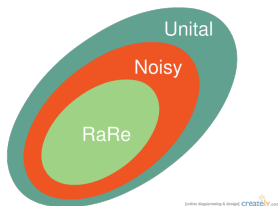
Channels \mathcal{N} of the form

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- Noisy operations are **unital**, but **not all** unital channels are noisy operations [Haagerup & Musat]!
- **RaRe channels** are noisy operations (via a tool in Lee & Selby's talk), but **not all** noisy operations are RaRe [Shor]!

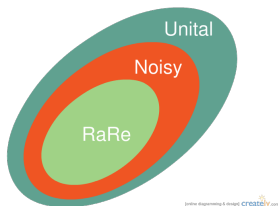
Conclusions

- We've analysed some different ways to define a resource theory of purity: **RaRe** channels, **noisy** operations, and **unital** channels.
















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- Despite the *strict* inclusions, they're equivalent: the induced preorder is the same (**majorisation**).

References

-  H Barnum, M P Müller, C Ududec, *New J. Phys.* **16** (12), 123029 (2014).
-  B Coecke, T Fritz, R W Spekkens, arXiv:1409.5531 [quant-ph] (2014).
-  G Chiribella, G M D'Ariano, P Perinotti, *Phys. Rev. A* **81** (6), 062348 (2010).
-  G Chiribella, CMS, *EPJ Web of Conferences* **95**, 03003 (2015).
-  G Chiribella, CMS, *New J. Phys.* **17** (10), 103027 (2015).
-  G Chiribella, CMS, *EPTCS* **195**, 96–115 (2015).
-  G Gour et al., *Phys. Rep.* **583**, 1–58 (2015).
-  U Haagerup, M Musat, *Commun. Math. Phys.* **303** (2), 555–594 (2011).
-  L Hardy, arXiv:1104.2066 [quant-ph] (2011).
-  M Horodecki, P Horodecki, J Oppenheim, *Phys. Rev. A* **67** (6), 062104 (2003).
-  P W Shor, *Structure of Unital Maps and the Asymptotic Quantum Birkhoff Conjecture*, presentation (2010).
-  L J Landau, R F Streater, *Linear Algebra Appl.* **193**, 107–127 (1993).
-  A Uhlmann, *Wiss. Z. Karl-Marx Univ. Leipzig* **20**, 633–637 (1971).